Thermal Modeling of the Solar Corona

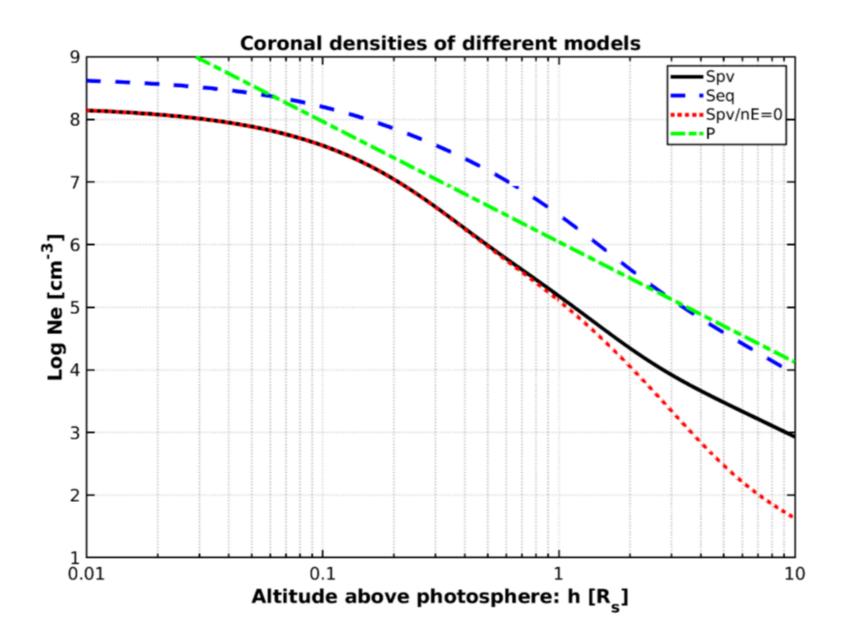
Thanassis Katsiyannis Royal Observatory of Belgium

Layout of the talk

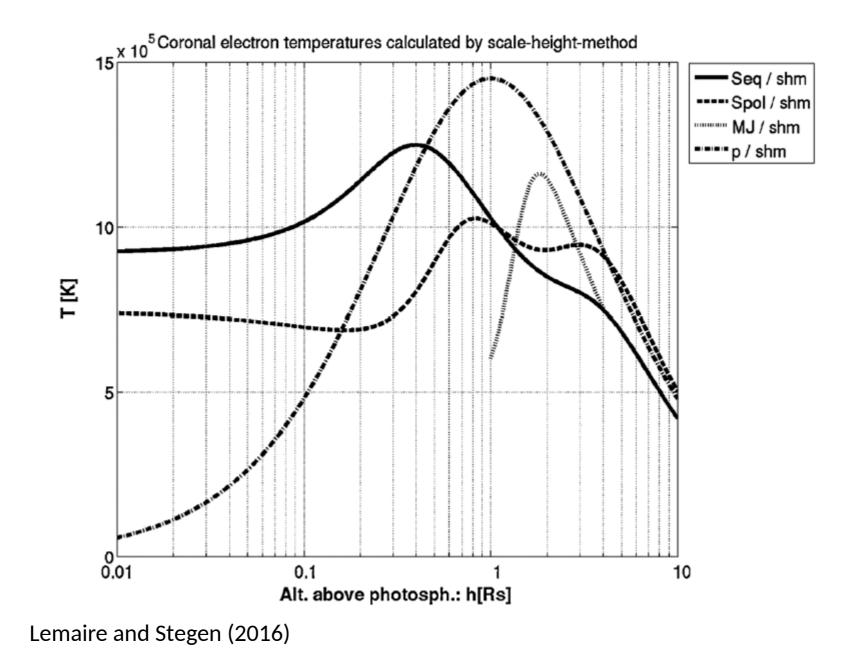
- Existing models of the ambient corona and their limitations
- The temperature profiles derived by the models

The Scale-Height model

- H = K*T / $\mu^* m_H^* g$ = [d(ln n_e)/dr]-1
- 1st problem: $g = g(r) = g_0^* R_{s^2} / r^2$ (Chapman 1957)
- 2nd problem: H not constant, i.e. corona not isothermal
- H not quasi-isothermal either (i.e. $d(\ln n_e)/dr$ not approximately exponential). Saito (1970)'s fit doesn't satisfy $p_e \approx n_e^* k * T$



Lemaire and Katsiyannis (2021)



The HST model

• Adds the dT/dr factor:

 $dT(r)/dr + [d(\ln n_e)/dr]^{-1} = -\mu^* m_H^* g^* R_{s^2}/(k r^2)$

- Alfven (1941) had already solved equations numerically.
- Problem: No solar wind!!

Parker's model

• Adds the p du/dt factor:

 $dT(r)/dr + [d(\ln n_e)/dr]^{-1} + \mu^* m_H^* g^* R_{S^2}/(k r^2) + \rho du/dt = 0,$

where $\rho = m_H n_e(r)$ and $n_e(r)$ decreases exponentially.

- 1st problem: Needs precise assumptions for the boundary condition of u₀ so u(r) is continuous and passes through a saddle point at the distance where the bulk speed becomes supersonic (Parker 1963).
- 2nd problem: It only diverges for T(∞)=0. Pottach (1960), Brandt et al (1965), and Gibson et al (1999) used T(∞)=0 and P(∞)=0. Even slightly different boundary conditions will produce divering steady state solutions.

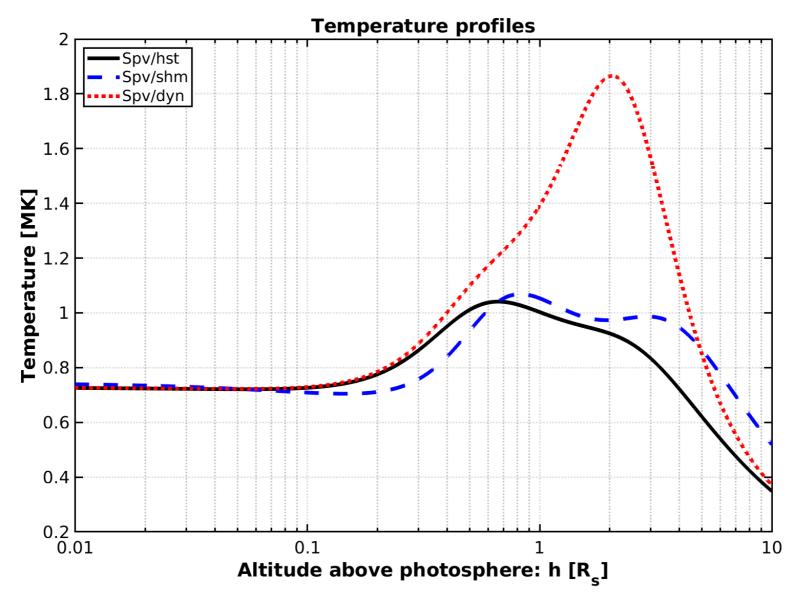
The DYN model

• Lemaire & Stegen (2016) also added at term to Saito (1970)'s fit. This was to correct for n_e(1AU). Saito (1970)'s fit up to 4 R_s:

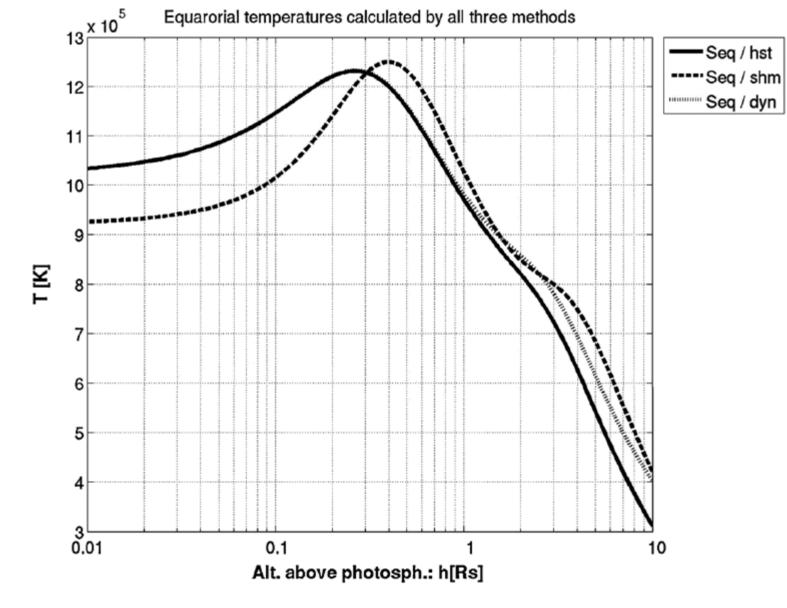
 $n_e(r)$ = 10⁸ [3.09 r⁻¹⁶ (1-0.5 sin(φ)) + 1.58 r⁻⁶ (1-0.95 sin(φ)) + 0.0251 r^{-2.5} (1-√sin(φ))] + $n_e(1AU)$ (215/r)²

- Boundary conditions (n_E and u_E) are set at 1 AU.
- $u(r,\phi) = u_E A_E / A(r) n_E / n_e(r,\phi)$,

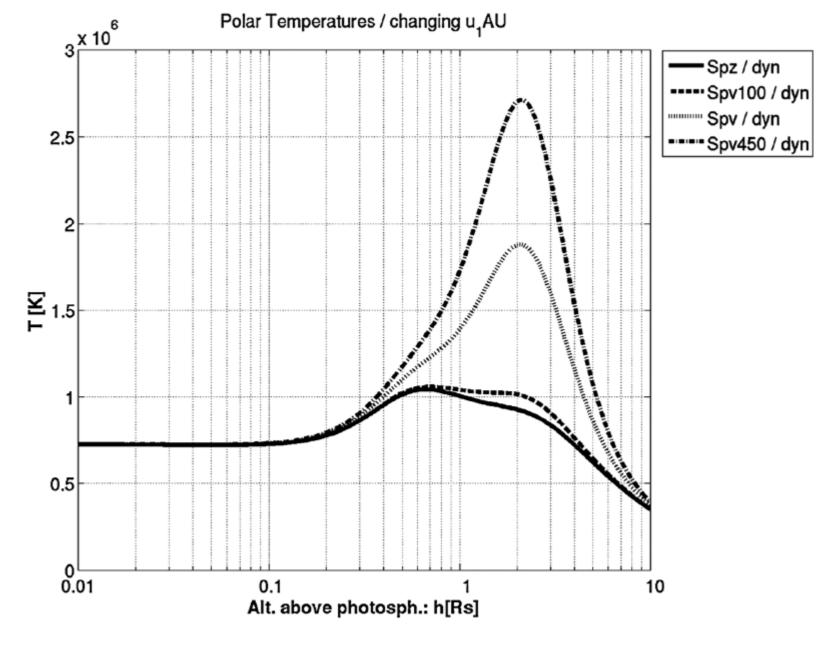
where A is the cross section of the flow tubes.



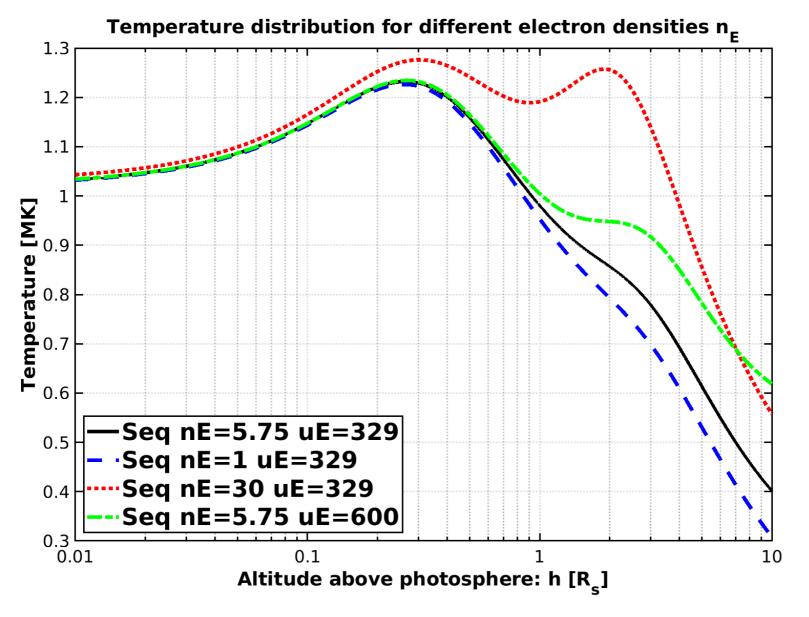
Lemaire and Katsiyannis (2021)



Lemaire and Stegen (2016)



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